Using LES/DNS Data for Neural Network-based Improvement of Existing Turbulence Models

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Turbulence Modeling Approaches

Traditional TM

Model the Reynolds stress (R_{ij})

- Boussinesq assumption, $R_{ij} \propto \text{velocity gradients}$
- Utilize turbulence physics to improve

Examples

- <u>Baldwin-Lomax</u> 0-eq
- Spallart-Allmaras 1-eq
 - Transport eq. for isotropic v_T
 - As many as 12 coefficients depending on PoV
 - Originally tuned to Samuel-Joubert flow and RAE 2822 airfoil
 - Many variations/improvements
- 2-eq/Algebraic R_{ij} models many more coefficients and variations

ML Enhanced TM

ML Assessment of Turbulence Physics

- Compute R_{ij} and turbulence source terms from LES/DNS data
- Determine how well existing turbulence models correlate to "truth" data
 - Directly assess the Boussinesq assumption
 - Directly assess derived turbulence physics

ML Optimization of TM model coefficients

- Experimental Objective Functions Yoder and Orkwis
- Utilize LES/DNS "truth" data
- Local or global variable driven

ML Classification of Flows

Ling and Templeton/Fuchi et al. variables

ML TM Development

 ML unsupervised learning to "discover" correlations between classification variables and equations based on them from LES/DNS data to derive new turbulence model forms

Reynolds Stresses

• "Reynolds stress," R_{ij} , is a consequence of the so-called Reynolds averaging process applied to the Navier-Stokes equations.

$$R_{ij} = -\rho \overline{u_i' u_j'}$$

 Derivatives of these terms appear in the Reynolds averaged Navier-Stokes equations (RANS).

$$\rho \bar{u}_j \frac{\partial \bar{u}_i}{\partial x_j} = \rho \bar{f}_i + \frac{\partial}{\partial x_j} \left[-\bar{p} \delta_{ij} + \mu \left(\frac{\partial \bar{u}_i}{\partial x_j} + \frac{\partial \bar{u}_j}{\partial x_i} \right) - \rho \bar{u}_i' \bar{u}_j' \right]$$
https://en.wikipedia.org/wiki/Reynolds-

averaged Navier%E2%80%93Stokes equations

Boussinesq proposed that

$$-\overline{v_i'v_j'} =
u_t \left(rac{\partial \overline{v_i}}{\partial x_j} + rac{\partial \overline{v_j}}{\partial x_i}
ight) - rac{2}{3}k\delta_{ij}$$

Which can be written in shorthand as

$$-\overline{v_i'v_j'}=2
u_tS_{ij}-rac{2}{3}k\delta_{ij}$$

where S_{ij} is the mean rate of strain tensor u_t is the turbulence eddy viscosity

$$k=rac{1}{2}\overline{v_i'v_i'}$$
 is the turbulence kinetic energy and δ_{ij} is the Kronecker delta.

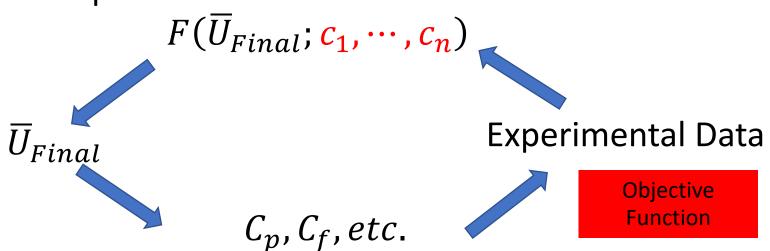
https://en.wikipedia.org/wiki/Turbulence modeling

 R_{ij} and subsequently ν_t or μ_t can be computed directly from LES/DNS data, embodying all the modeling assumptions made in the scheme.

Turbulence Modeling

Turbulence Model Process

$$\overline{U} \to F(\overline{U}) \to \mu_T | Re_{\sigma} \to TST \qquad \qquad \overline{U}_{Fina}$$



Algebraic Turbulence Models

 Johnson and King, Cebici and Smith, and Baldwin and Lomax invented models using Prandtl's mixing length idea to form algebraic turbulence models that work well near walls.

$$\mu_t = \begin{cases} \mu_{tinner} & \text{if } y \leq y_{crossover} \\ \mu_{touter} & \text{if } y > y_{crossover} \end{cases}$$

Where $y_{crossover}$ is the smallest distance from the surface where μ_{tinner} is equal to μ_{touter}

$$y_{crossover} = MIN(y) : \mu_{tinner} = \mu_{touter}$$

The inner region is given by the Prandtl - Van Driest formula:

$$\mu_{tinner} = \rho l^2 |\Omega|$$

$$|\Omega| = \sqrt{2\Omega_{ij}\Omega_{ij}}$$

$$\Omega_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial x_i} \right)$$

The outer region is given by:

$$\mu_{touter} = \rho K C_{CP} F_{WAKE} F_{KLEB}(y)$$

Where

$$F_{WAKE} = MIN \left(y_{MAX} F_{MAX} ; C_{WK} y_{MAX} \frac{u_{DIF}^2}{F_{MAX}} \right)$$

 y_{MAX} and F_{MAX} are determined from the maximum of the function:

$$F(y) = y |\Omega| \left(1 - e^{\frac{-y^{+}}{A^{+}}}\right)$$

 F_{KLER} is the intermittency factor given by:

$$F_{KLEB}(y) = \left[1 + 5.5 \left(\frac{y C_{KLEB}}{y_{MAX}}\right)^{6}\right]^{-1}$$

Introducing the essential idea of fitting coefficients to match experiments.

Notably:

Garabedian-Korn airfoil Horstmann-Hung ramp Hopkins-Inouye plate

 A^{+} C_{CP} C_{KLEB} C_{WK} k K26 1.6 0.3 0.25 0.4 0.0168

 u_{DIF} is the difference between maximum and minimum speed in the profile. For boundary layers the minimum is always set to zero

$$u_{DIF} = MAX(\sqrt{u_i u_i}) - MIN(\sqrt{u_i u_i})$$

Spalart-Allmaras 1-eq Model

• Spalart and Allmaras created a 1-equation model for the transport of ν_t

$$\begin{split} \nu_t &= \tilde{\nu} f_{v1}, \quad f_{v1} = \frac{\chi^3}{\chi^3 + C_{v1}^3}, \quad \chi := \frac{\tilde{\nu}}{\nu} \\ \frac{\partial \tilde{\nu}}{\partial t} + u_j \frac{\partial \tilde{\nu}}{\partial x_j} &= C_{b1} [1 - f_{t2}] \tilde{S} \tilde{\nu} + \frac{1}{\sigma} \{ \nabla \cdot [(\nu + \tilde{\nu}) \nabla \tilde{\nu}] + C_{b2} |\nabla \tilde{\nu}|^2 \} - [C_{w1} f_w - \frac{C_{b1}}{\kappa^2} f_{t2}] \left(\frac{\tilde{\nu}}{d} \right)^2 + f_{t1} \Delta U^2 \\ \tilde{S} &\equiv S + \frac{\tilde{\nu}}{\kappa^2 d^2} f_{v2}, \quad f_{v2} = 1 - \frac{\chi}{1 + \chi f_{v1}} \\ \text{where} \\ S &= \equiv \sqrt{2 \Omega_{ij} \Omega_{ij}} \\ \Omega_{ij} &\equiv \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial x_i} \right) \\ f_w &= g \left[\frac{1 + C_{w3}^6}{g^6 + C_{w3}^6} \right]^{1/6}, \quad g = r + C_{w2} (r^6 - r), \quad r \equiv \frac{\tilde{\nu}}{\tilde{S} \kappa^2 d^2} \\ f_{t1} &= C_{t1} g_t \exp \left(-C_{t2} \frac{\omega_t^2}{\Delta U^2} [d^2 + g_t^2 d_t^2] \right) \\ f_{t2} &= C_{t3} \exp(-C_{t4} \chi^2) \\ \text{d is the distance to the closest surface} \end{split}$$

The constants are

$$\sigma = 2/3$$
 $C_{b1} = 0.1355$
 $C_{b2} = 0.622$
 $\kappa = 0.41$
 $C_{w1} = C_{b1}/\kappa^2 + (1 + C_{b2})/\sigma$
 $C_{w2} = 0.3$
 $C_{w3} = 2$
 $C_{v1} = 7.1$
 $C_{t1} = 1$
 $C_{t2} = 2$
 $C_{t3} = 1.1$
 $C_{t4} = 2$

Up to 12 depending on how you count them.
Their case was correlated against Samuel-Joubert and RAE 2822 airfoil

Reynolds Stress Models

Transport equations for R_{ij} are derived in Reynolds Stress Transport models

$$\frac{\partial}{\partial t} \left(-\tau_{ij}^{\mathsf{T}} \right) + \frac{\partial}{\partial x_{k}} \left(-\tau_{ij}^{\mathsf{T}} \tilde{u}_{k} \right) = \mathcal{P}_{ij} + \mathcal{D}_{ij} + \Pi_{ij} - \frac{2}{3} \bar{\rho} \epsilon \delta_{ij}$$

where the production, diffusion, and pressure-strain terms are

$$\mathcal{P}_{ij} = \tau_{ik}^{\scriptscriptstyle T} \frac{\partial \tilde{u}_j}{\partial x_k} + \tau_{jk}^{\scriptscriptstyle T} \frac{\partial \tilde{u}_i}{\partial x_k}$$

$$\mathcal{D}_{ij} = \frac{\partial}{\partial x_k} \left[\left(\mu_{\scriptscriptstyle L} + \frac{\mu_{\scriptscriptstyle T}}{\sigma_k} \right) \frac{\partial}{\partial x_k} \left(\frac{-\tau_{ij}^{\scriptscriptstyle T}}{\bar{\rho}} \right) \right]$$

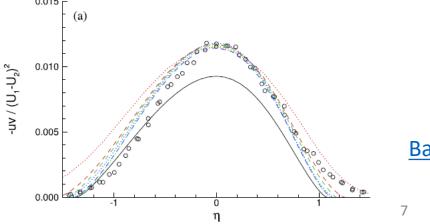
$$\Pi_{ij} = -\left(C_1^0 + C_1^1 \frac{\mathcal{P}}{\bar{\rho}\epsilon}\right) \bar{\rho}\epsilon b_{ij} + C_2 \bar{\rho}k \widetilde{S}_{ij}^D
+ C_3 \bar{\rho}k \left[b_{ik}\widetilde{S}_{kj} + \widetilde{S}_{ik}b_{kj} - \frac{2}{3}b_{mn}\widetilde{S}_{mn}\delta_{ij}\right] - C_4 \bar{\rho}k \left[b_{ik}\widetilde{R}_{kj} - \widetilde{R}_{ik}b_{kj}\right]$$

Which produce more and more coefficients.

An attempt at such modeling was made by Yoder and Orkwis using ML ideas against Delville's jet data.

Yoder, D. A., & Orkwis, P. D. (2021). On the use of optimization techniques for turbulence model calibration. *Computers & Fluids*, 214, 104752.

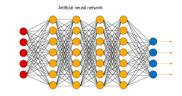




Back

ML Turbulence Modeling - 1

Turbulence Model Process



$$\overline{U} \to F(\overline{U}) \to \mu_T | Re_{\sigma} \to TST$$



 \overline{U}_{Final}

$$F(\overline{U}_{Final}; c_1, \cdots, c_n)$$
 \overline{U}_{Final}
Experimental Data
 $C_p, C_f, etc.$
Objective
Function

NN Model Accuracy

ML Model

TensorFlow NN used with a tanh loss function at all layers

12 node input layer

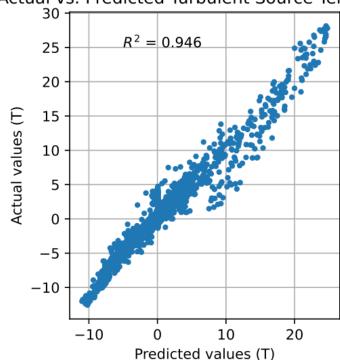
3 hidden layers of 30 nodes

1 output

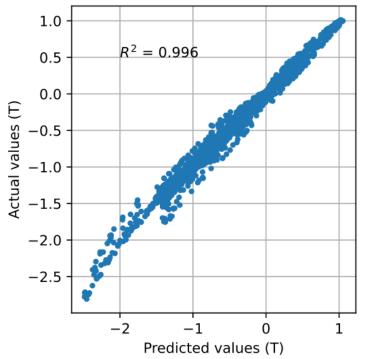
To show validity of using a NN to fit lower fidelity data, used an 80-20 train-test split

Trained the X and Y turbulent source terms independently

Actual vs. Predicted Turbulent Source Terms (NN)







X Turbulent Source Term

Turbulent Source Term

The inferred turbulent source term can be computed from the LES/DNS data. This is the term that should be compared to the turbulence model source term.

$$\rho \overline{u_j} \frac{\partial \overline{u_i}}{\partial x_j} = \rho \overline{f_i} + \frac{\partial}{\partial x_j} \left[-\overline{p} \delta_{ij} + \mu \left(\frac{\partial \overline{u_i}}{\partial x_j} + \frac{\partial \overline{u_j}}{\partial x_i} \right) - \rho \overline{u_i' u_j'} \right]$$

So that the governing equations become:

$$\rho \overline{u_j} \frac{\partial \overline{u_i}}{\partial x_j} = \rho \overline{f_i} + \frac{\partial}{\partial x_j} \left[-\overline{p} \delta_{ij} + \mu \left(\frac{\partial \overline{u_i}}{\partial x_j} + \frac{\partial \overline{u_j}}{\partial x_i} \right) \right] + \mathbf{T_i}$$

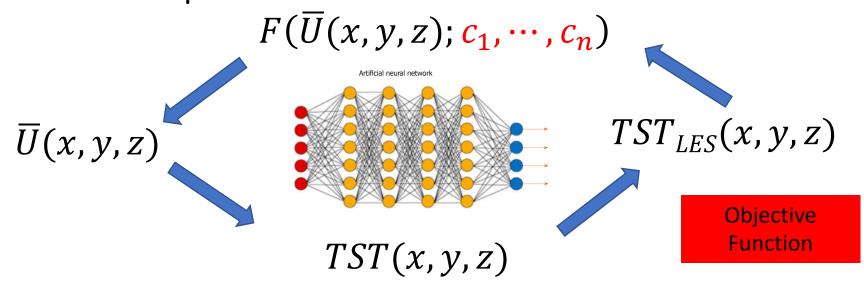
Where the LES/DNS data provides: $T_i = -\frac{\partial \rho u_i' u_j'}{\partial x_i}$

And turbulence modeling uses: $\widetilde{T}_i = \frac{\partial}{\partial x_j} \left[\mu_t \left(\frac{\partial \overline{u_i}}{\partial x_j} + \frac{\partial \overline{u_j}}{\partial x_i} \right) \right]$ or some variation.

ML Turbulence Modeling - 2

Turbulence Model Process

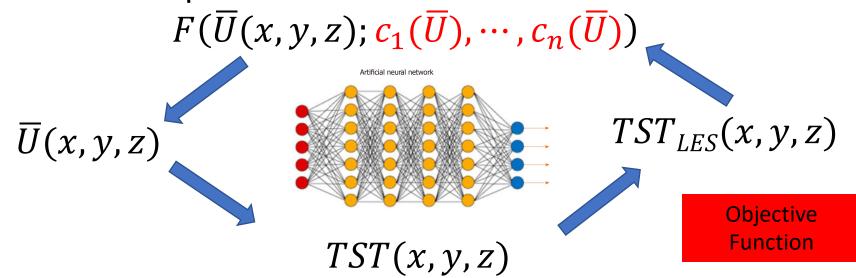
$$\overline{U} \to F(\overline{U}) \to \mu_T | Re_{\sigma} \to TST \qquad \qquad \overline{U}_{Final}$$



ML Turbulence Modeling - 3

Turbulence Model Process

$$\overline{U} \to F(\overline{U}) \to \mu_T | Re_{\sigma} \to TST \qquad \qquad \overline{U}_{Final}$$



Classification: Ling and Templeton Variables

085103-11 J. Ling and J. Templeton

Phys. Fluids 27, 085103 (2015)

TABLE II. Non-dimensional inputs.

#	Description	Formula	#	Description	Formula
1	Non-dimensionalized Q criterion	$\frac{\ R\ ^2 - \ S\ ^2}{\ R\ ^2 + \ S\ ^2}$	7	Ratio of pressure normal stresses to normal shear stresses	$\frac{\sqrt{\frac{\partial P}{\partial x_i}} \frac{\partial P}{\partial x_i}}{\sqrt{\frac{\partial P}{\partial x_f} \frac{\partial P}{\partial x_f}} + 0.5\rho \frac{\partial U_k^2}{\partial x_k}}$
2	Turbulence intensity	$\frac{k}{0.5U_iU_i+k}$	8	Vortex stretching	$\frac{\sqrt{\omega_{l}\frac{\partial U_{i}}{\partial x_{l}}\omega_{k}\frac{\partial U_{i}}{\partial x_{k}}}}{\sqrt{\omega_{l}\frac{\partial U_{n}}{\partial x_{l}}\omega_{m}\frac{\partial U_{n}}{\partial x_{m}}} + S }}{ U_{k}U_{l}\frac{\partial U_{k}}{\partial x_{l}} }$
3	Turbulence Reynolds number	$\min(\frac{\sqrt{k}\ d}{50\nu},2)$	9	Marker of Gorle et al., ¹³ deviation from parallel shear flow	$\frac{ U_k U_l \frac{\partial U_k}{\partial x_l} }{\sqrt{U_n U_n U_i \frac{\partial U_i}{\partial x_f} U_m \frac{\partial U_m}{\partial x_f}}} + U_i U_f \frac{\partial U_i}{\partial x_f}$
4	Pressure gradient along streamline	$\frac{U_k \frac{\partial P}{\partial x_k}}{\sqrt{\frac{\partial P}{\partial x_f} \frac{\partial P}{\partial x_f} U_i U_i} + U_l \frac{\partial P}{\partial x_l} }$	10	Ratio of convection to production of <i>k</i>	$\frac{U_i \frac{dk}{\partial x_i}}{ u'_f u'_l S_{fl} + U_l \frac{dk}{\partial x_l}}$
5	Ratio of turbulent time scale to mean strain time scale	$\frac{\ S\ k}{\ S\ k+\epsilon}$	11	Ratio of total Reynolds stresses to normal	$\frac{\ \overline{u_i'u_j'}\ }{k+\ \overline{u_i'u_j'}\ }$
6	Viscosity ratio	$\frac{v_t}{100v + v_t}$	12	Reynolds stresses Cubic eddy viscosity comparison	$\frac{S_{if}(\overline{u_i'u_f'}_{\text{CEVM}} - \overline{u_i'u_f'}_{\text{LEVM}})}{S_{kl}(\overline{u_k'u_l'}_{\text{CEVM}} + \overline{u_k'u_l'}_{\text{LEVM}})}$

What's Needed...

- Clean LES/DNS data that has been run sufficiently long for meaningful statistical data.
 - High quality grids
 - Sufficient T_c
 - Mean flow solution throughout the field
 - Reynolds stresses throughout the field
 - Able to create TSTs at each point
- A good idea of what variables best characterize a flow.
 - Including the scale at which to characterize (global, local, cell).
- A broad range of "different" solutions.

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Questions?